



# Theory of the hadron (baryon) spectrum

- Missing resonances,  $N^*$  program at Jlab
- Unitary, multi-channel analysis
- Competing models of the short-distance  $qq$  interaction
  - Can this issue be resolved by examining mixing angles in  $S_{11}$  and  $D_{13}$  states?
- New insights from lattice QCD calculations
  - Flux tube potential picture of  $Q\bar{Q}$ ,  $QQQ$
  - String breaking
  - Dramatic behavior of baryon masses in low  $m_\pi$  limit



# Missing resonances

- Symmetric quark models predict many positive (and doubly-excited negative) parity states not seen in analyses of data
- PDG states were based on analyses of  $N\pi$  elastic scattering
  - Are there states which couple weakly to  $N\pi$ ?
  - These can be expected to be “missing”
  - **Evidence** for them should show up in other ( $N\pi\pi$ ,  $\Lambda K$ ,...) final states, excited with EM probes from nucleon targets (make  $N^*$  or  $\Delta^*$ )
  - Their existence will be **established** in multi-channel analyses of many different final states

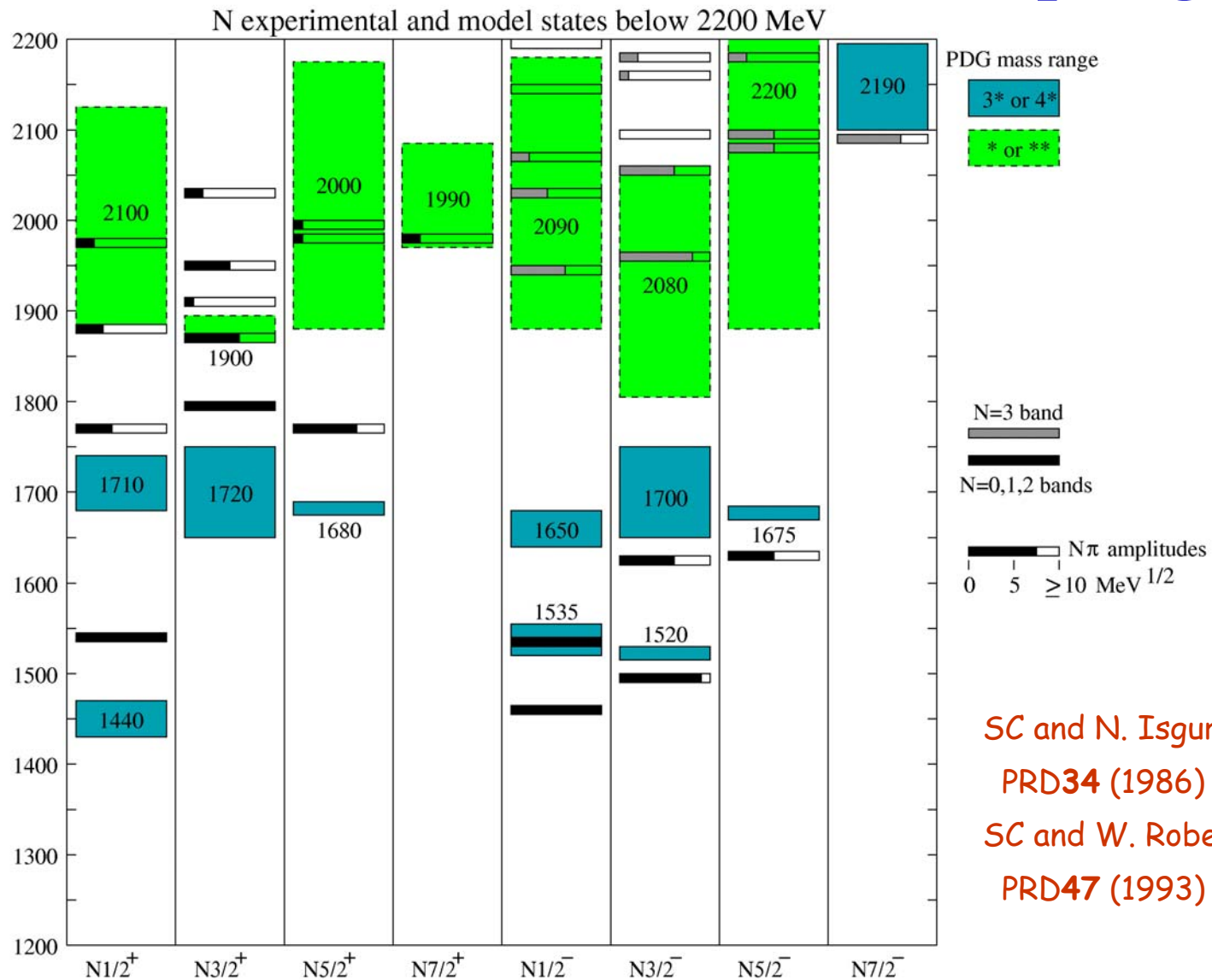


# Missing resonances...

- R. Koniuk and N. Isgur
  - Elementary-meson emission decay model (quarks directly emit mesons)
  - Baryon structure from Isgur-Karl model
  - ⇒ Showed missing states correspond in number and energy to those predicted to decouple from  $N\pi$
- $^3P_0$ : quark-pair creation decay model
  - Hadrons decay by creation of quark pair with quantum numbers of the vacuum
    - String breaking has same structure
  - Emitted mesons have structure
  - Baryons wave functions from structure model
  - ⇒ Can correlate many decays with a couple of parameters, predict which final states will show evidence of new states



# Nucleon model states and $N\pi$ couplings

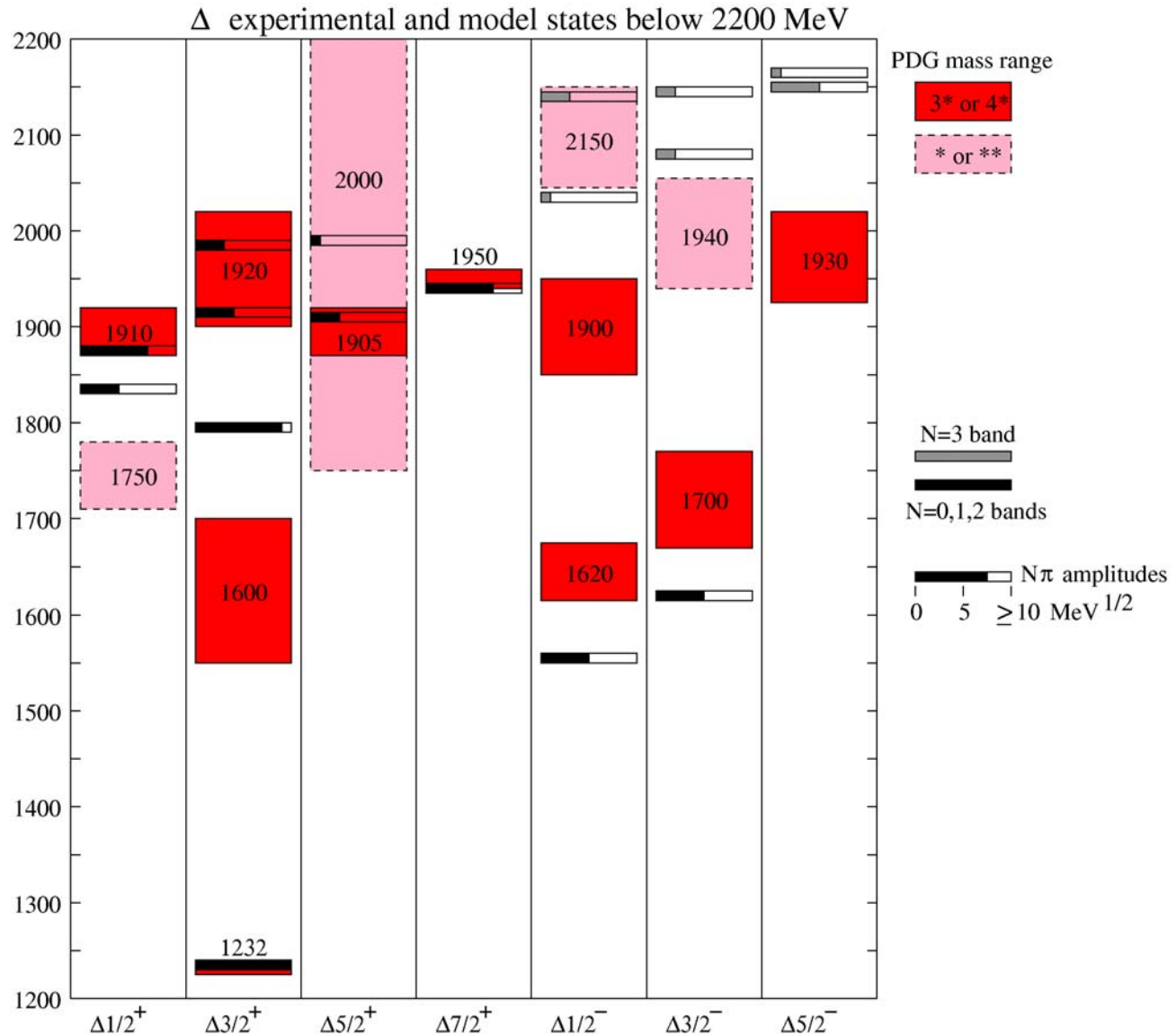


SC and N. Isgur,  
PRD**34** (1986) 2809;  
SC and W. Roberts,  
PRD**47** (1993) 2004





# $\Delta$ model states and $N\pi$ couplings

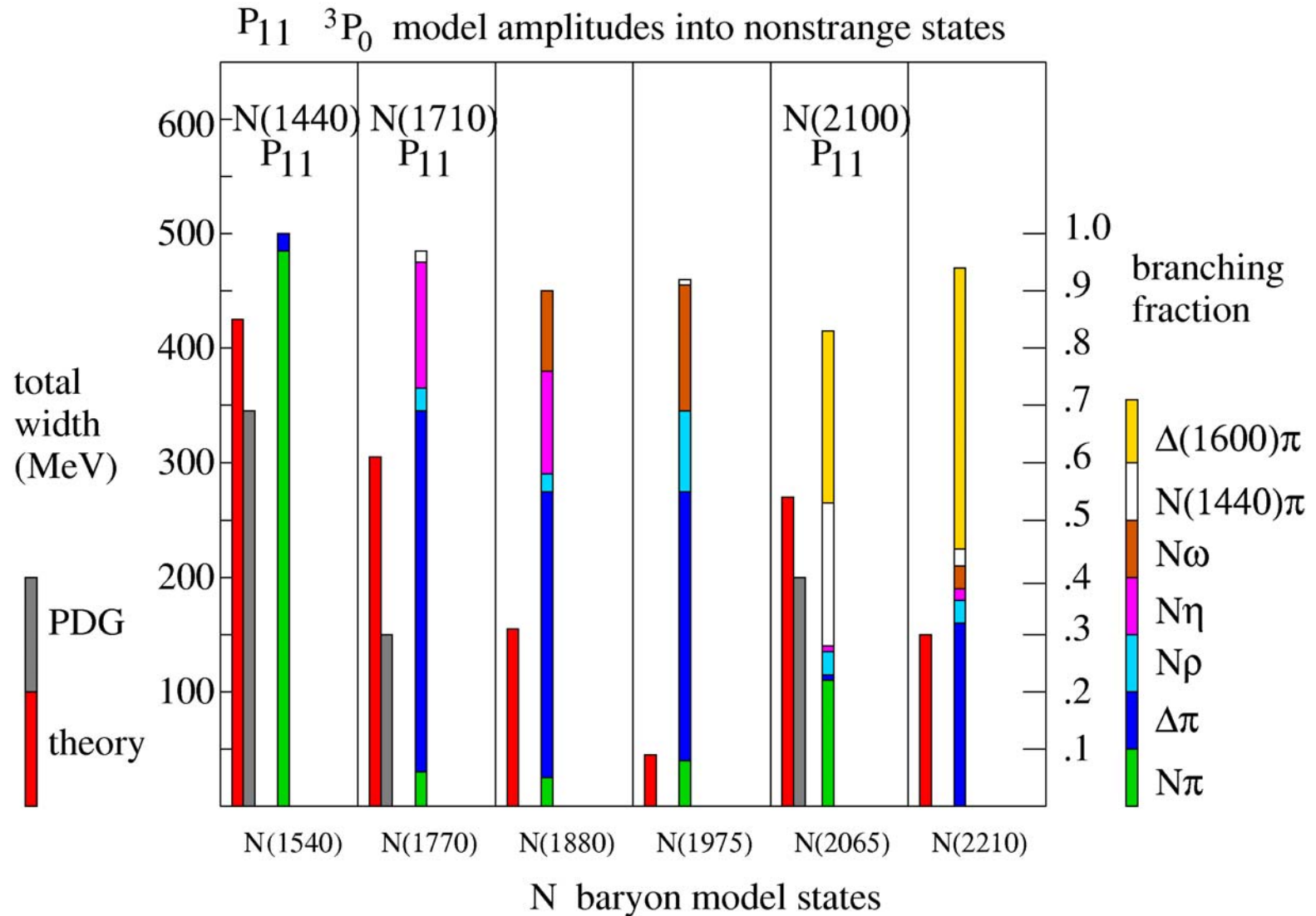


# Predictions for branching fractions

- Can calculate decay branching fractions if calculate enough decay channels!
  - Have looked at decays of  $N^*$ ,  $\Delta^*$  states to  $N$ ,  $\Delta$ ,  $N(1440)$ ,  $\Delta(1600)$ ... and  $\pi$ ,  $\eta$ ,  $\eta'$ ,  $\rho$ ,  $\omega$
  - Also decays to strange final states, including  $\Lambda$ ,  $\Sigma$ ,  $\Sigma^*$ ,  $\Lambda(1405)$ ,... and  $K$ ,  $K^*$  (not shown below)



# Predictions for branching fractions...



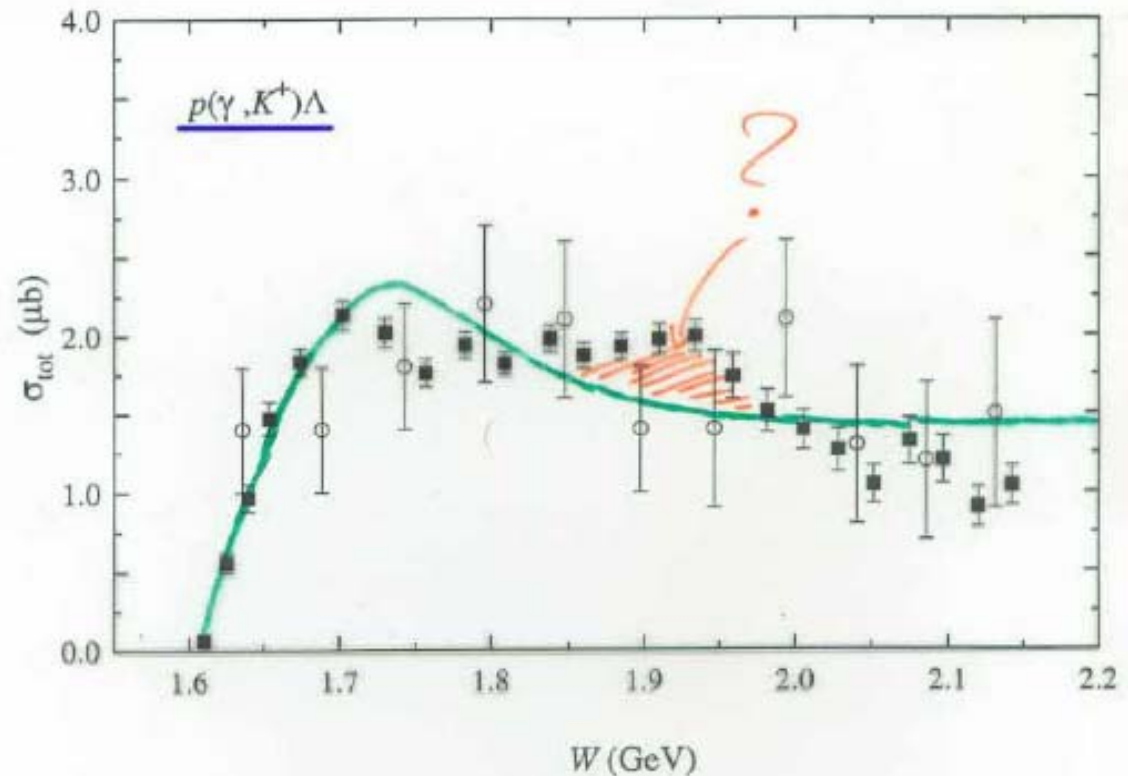


# Have we seen new resonances?

- Effect seen in  $\gamma p$  to  $\Lambda K^+$  at  $W \sim 1900$  MeV at SAPHIR; **C. Bennhold et al.**

- Could be  $D_{13}$  predicted by model (has sizeable  $\gamma N$  and  $\Lambda K$  coupling)
- Or the effect of more than one state

- Need coupled channels to confirm



# Have we seen new resonances?...

- Possible third  $S_{11}$  resonance at  $\sim 1.8$  GeV
  - Rapidly rising cross section for  $\gamma N$  to  $N\eta$  above existing  $S_{11}$  states at GRAAL
  - Similar rise seen in CLAS in  $\eta$  photoproduction experiments
  - Very interesting possibility for quark models!
- Evidence from  $N\rho$ ,  $\Delta\pi$  electroproduction in CLAS that  $P_{13}(1720)$  region not understood
  - Two states? Hard to fit in models...
- Huge potential for interesting new results from Jlab experiments!
  - Need multi-channel analysis to firmly establish new states



# Unitary multi-channel analysis

- Based on K-matrix: KSU group, D.M. Manley and collaborators
  - Effective Lagrangians: T. Sato and T.-S. H. Lee; GWU group: C. Bennhold, H. Haberzettl,...
  - CMU-LBL (Cutkosky) method: Vrana, Dytman and Lee, Phys. Rept. 328, (200), 181 (Pitt-ANL)
    - Extend  $\pi N \Rightarrow \pi N$  analysis to many other channels
    - $N\pi$ ,  $N\eta$ ,  $N\gamma$ ,  $N\rho$ ,  $\Delta\pi$ ,  $N\sigma$ ,  $\Lambda K$ ,  $\Sigma K$ ,...
    - Incorporate constraints from analyticity
    - Impose multi-channel unitarity
- ⇒ Resonance parameters are extracted from analysis of partial-wave amplitudes in all open channels simultaneously



## Coupled Channel Picture of Resonance Excitation

- Each resonance can be reached through each asymptotic channel
- T matrix provides unitary, analytic structure
- all channels (e.g.  $\pi N$ ,  $\rho N$ ) couple to all other channels in intermediate state
- photon multipoles ( $E_{l\pm}$ ,  $M_{l\pm}$ ) directly related to T

$$T(J^\pi) = \begin{pmatrix} T_{\pi N \rightarrow \pi N} & T_{\eta N \rightarrow \pi N} & T_{\gamma N \rightarrow \pi N} & T_{\rho N \rightarrow \pi N} & T_{\sigma N \rightarrow \pi N} & T_{K\Lambda \rightarrow \pi N} & T_{K\Sigma \rightarrow \pi N} \\ T_{\pi N \rightarrow \eta N} & T_{\eta N \rightarrow \eta N} & T_{\gamma N \rightarrow \eta N} & T_{\rho N \rightarrow \eta N} & T_{\sigma N \rightarrow \eta N} & T_{K\Lambda \rightarrow \eta N} & T_{K\Sigma \rightarrow \eta N} \\ T_{\pi N \rightarrow \gamma N} & T_{\eta N \rightarrow \gamma N} & T_{\gamma N \rightarrow \gamma N} & T_{\rho N \rightarrow \gamma N} & T_{\sigma N \rightarrow \gamma N} & T_{K\Lambda \rightarrow \gamma N} & T_{K\Sigma \rightarrow \gamma N} \\ T_{\pi N \rightarrow \rho N} & T_{\eta N \rightarrow \rho N} & T_{\gamma N \rightarrow \rho N} & T_{\rho N \rightarrow \rho N} & T_{\sigma N \rightarrow \rho N} & T_{K\Lambda \rightarrow \rho N} & T_{K\Sigma \rightarrow \rho N} \\ T_{\pi N \rightarrow \sigma N} & T_{\eta N \rightarrow \sigma N} & T_{\gamma N \rightarrow \sigma N} & T_{\rho N \rightarrow \sigma N} & T_{\sigma N \rightarrow \sigma N} & T_{K\Lambda \rightarrow \sigma N} & T_{K\Sigma \rightarrow \sigma N} \\ T_{\pi N \rightarrow K\Lambda} & T_{\eta N \rightarrow K\Lambda} & T_{\gamma N \rightarrow K\Lambda} & T_{\rho N \rightarrow K\Lambda} & T_{\sigma N \rightarrow K\Lambda} & T_{K\Lambda \rightarrow K\Lambda} & T_{K\Sigma \rightarrow K\Lambda} \\ T_{\pi N \rightarrow K\Sigma} & T_{\eta N \rightarrow K\Sigma} & T_{\gamma N \rightarrow K\Sigma} & T_{\rho N \rightarrow K\Sigma} & T_{\sigma N \rightarrow K\Lambda} & T_{K\Lambda \rightarrow K\Sigma} & T_{K\Sigma \rightarrow K\Sigma} \end{pmatrix}$$



# Results of Pitt-ANL analysis

state	Mass (MeV)	Width (MeV)	$B_{N\pi}(\%)$	$B_{N\eta}(\%)$
$S_{11}(1535)$ (PDG)	1545(3) 1520-1555	127(19) 100-250	35(4) 35-55	55(5) 30-55
$S_{11}(1650)$ (PDG)	1693(12) 1640-1680	225(40) 145-190	73(2) 55-90	-2(1) 3-10
$D_{13}(1520)$ (PDG)	1520(3) 1515-1530	118(4) 110-135	61(2) 50-60	0(1)
$D_{13}(1700)$ (PDG)	1729(33) 1650-1750	178(133) 50-150	4(1) 5-15	7(1)



# QCD-inspired models of baryons

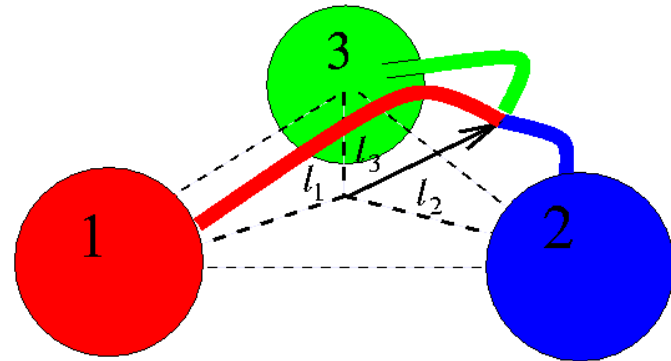
- Work in adiabatic (Born-Oppenheimer) approximation
- Effect of glue is to generate  $qq$  potentials
  - Confining potential:
    - flux-tube?
    - Dirac structure in relativistic treatments?
  - Short-distance (residual) interactions
    - OGE model: link to meson and heavy-quark physics
    - OBE model: chiral symmetry dominates soft QCD spectrum
    - Instanton-induced model: QCD vacuum structure implies 't Hooft's instanton-induced interaction





# Effective degrees of freedom

- One popular picture:
  - Constituent quarks
  - Glue in flux tubes (confinement)



⇒ Constituent quarks:

⇒ Dynamically generated constituent masses, which run with  $Q^2$  :

$$m_{u,d} \approx K_{u,d} \approx \Lambda_{\text{QCD}} \approx 200 \text{ MeV}$$

⇒ Effective sizes, form factors



# What are their residual interactions?

- Ground-state spectrum suggests flavor-dependent short-range (contact) interactions
- One-gluon exchange: good fit to ground states with (color-magnetic dipole-dipole), e.g.  $\Delta$ -N,  $\Sigma$ - $\Lambda$   
 $\Rightarrow$  DeRujula, Georgi, Glashow

$$M = \sum_{i=1}^3 m_i + \frac{2\alpha_s}{3} \frac{8\pi}{3} \langle \delta^3(\mathbf{r}) \rangle \sum_{i<j=1}^3 \frac{\mathbf{S}_i \cdot \mathbf{S}_j}{m_i m_j}$$

- Explains regularities in meson spectrum (e.g. evolution of vector-pseudoscalar splitting with quark mass)
  - Unclear why this should work for light quarks...
- Taken at face value predicts tensor interaction

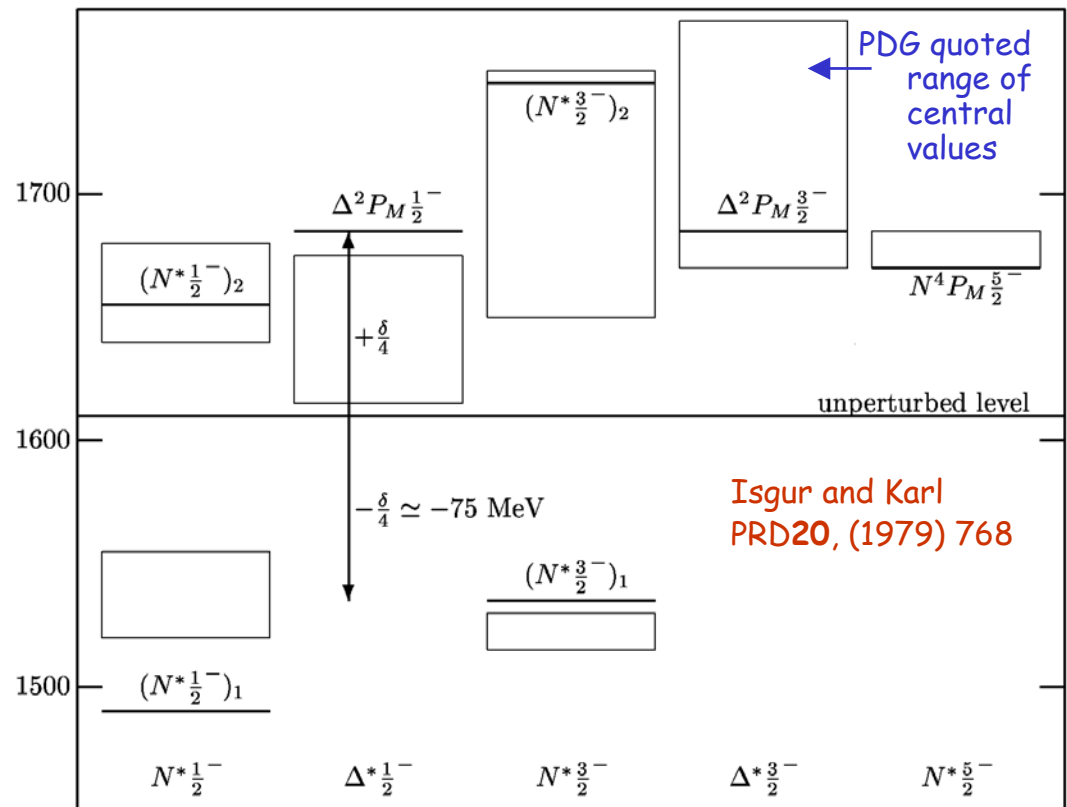
$$H_{\text{hyp}}^{ij} = \frac{2\alpha_s}{3m_i m_j} \left\{ \frac{8\pi}{3} \mathbf{S}_i \cdot \mathbf{S}_j \delta^3(\mathbf{r}_{ij}) + \frac{1}{r_{ij}^3} \left[ \frac{3(\mathbf{S}_i \cdot \mathbf{r}_{ij})(\mathbf{S}_j \cdot \mathbf{r}_{ij})}{r_{ij}^2} - \mathbf{S}_i \cdot \mathbf{S}_j \right] \right\}$$

- And spin-orbit interactions, at a level not present in spectrum



# Residual interactions...

- Isgur and Karl PRD20, (1979) 768
- Contact splitting active in L=1 excited states
- Characteristic splitting is  $(m_{\Delta} - m_N)/2$
- Add consistent tensor interaction
- No strong evidence for tensor from spectrum
- Best evidence from decays,  $S_{11}(1535) \rightarrow N\eta$



# Residual interactions...

- Variational calculation in large HO basis (SC, N. Isgur)
  - String confinement, plus associated spin-orbit
  - Include OGE Coulomb, contact, tensor, spin-orbit
  - Relativistic KE, extended constituent quarks
  - parameterize momentum dependence, running coupling...

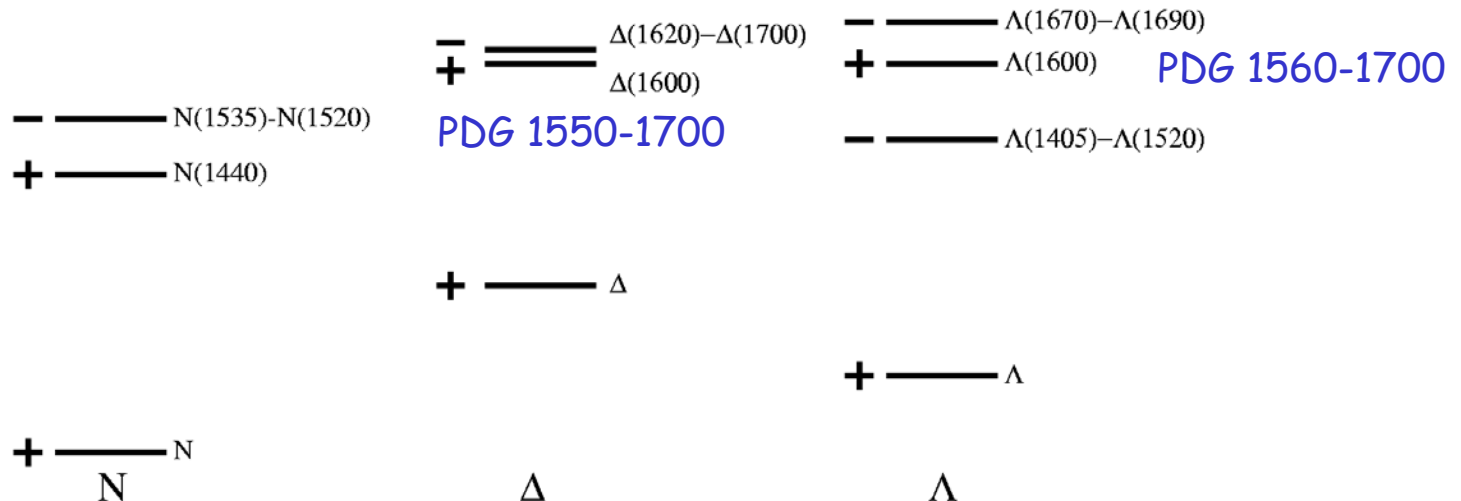
$$\left(\frac{m_i m_j}{E_i E_j}\right)^{\frac{1}{2} + \epsilon_{\text{cont}}} \frac{8\pi}{3} \alpha_s(r_{ij}) \frac{2}{3} \frac{\mathbf{S}_i \cdot \mathbf{S}_j}{m_i m_j} \left[ \frac{\sigma_{ij}^3}{\pi^{\frac{3}{2}}} e^{-\sigma_{ij}^2 r_{ij}^2} \right] \left(\frac{m_i m_j}{E_i E_j}\right)^{\frac{1}{2} + \epsilon_{\text{cont}}}$$

- Photocouplings calculated with  $H_{\text{int}}$  expanded to  $O(p^2/m^2)$
  - Strong decays calculated in pair creation ( ${}^3P_0$ ) model (with W. Roberts)
- Reasonable agreement; allows prediction of favorable channels to find 'missing' baryons
- Puzzles: Roper mass;  $\Lambda_{3/2}$ -(1520)- $\Lambda_{1/2}$ -(1405); L=1 too light by 50 MeV, positive parity too massive by 50 MeV,...



# Residual interactions...

- Another possibility: should light quarks exchange pions?  
Robson; Buchmann, Faessler,...
- Claim gluons not active in light-quark hadrons: flavor dependence through exchange of octet of pseudoscalars (GBE)
- Contact interaction: 
$$H_\chi \sim - \sum_{i < j} \frac{V(\mathbf{r}_{ij})}{m_i m_j} \lambda_i^F \cdot \lambda_j^F \sigma_i \cdot \sigma_j$$
- Order of states inverted? Natural with GBE  
⇒ Glozman & Riska (GR)



# Residual interactions...

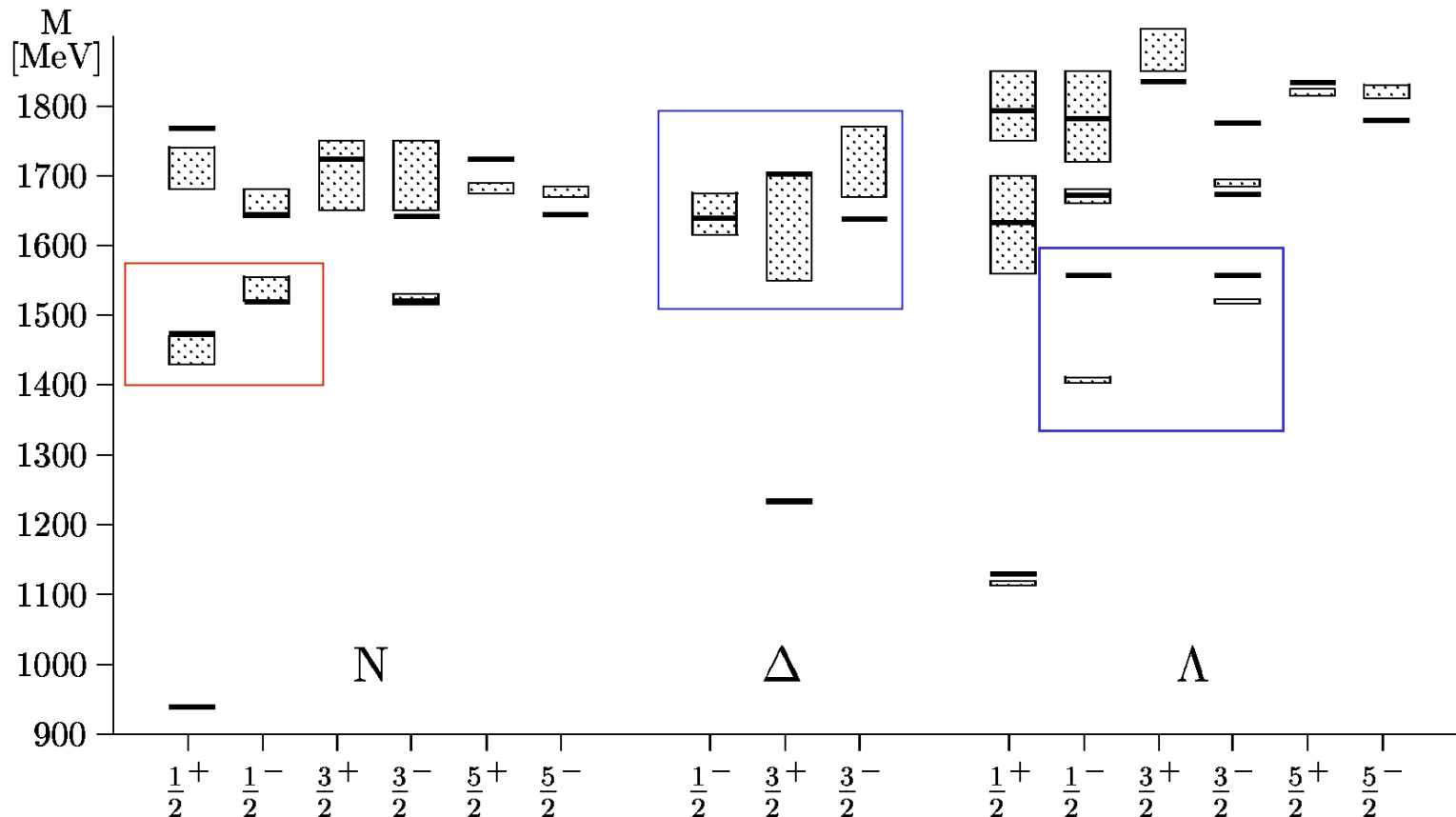
- GR fit radial matrix elements of  $V(r_{ij})$  to spectrum
- **Calculated** in variational H.O. basis with consistent tensor **Glozman, Plessas, Theussl, Wagenbrunn, & Varga**
- Add nonets of exchange vector mesons and scalars
  - Relativistic K.E., string confinement; calculate decays
  - Problems can arise in some of these models with decays (poor results with OBE)
  - E.g.  $N\eta$  decay of  $S_{11}(1535)$  not possible without vector exchange added to GBE





# Residual interactions...

- Results for spectrum:



# Residual interactions...

- Another flavor-dependent possibility: instanton-induced interactions
- Present if qq in S-wave,  $I=0$ ,  $S=0$  state

$$\langle q^2; S, L, T | W | q^2; S, L, T \rangle = -4g \delta_{S,0} \delta_{L,0} \delta_{T,0} \mathcal{W}$$

- $W$  is a contact interaction (has range  $\lambda$ )
  - causes no shifts in  $\Delta^*$  masses
  - No tensor interaction, or spin-orbit forces
- Applied to excited states **Blask, Bohn, Huber, Metsch & Petry**
  - solve Bethe-Salpeter equation



# Instanton-induced interactions

- Quarks confined by linear q-q potential

$$V(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = A_3 + B_3 \sum_{i < j} |\mathbf{r}_i - \mathbf{r}_j|$$

- Relativistic treatment, so need to choose Dirac structure of potential

$$A_3 = a \frac{3}{4} \left[ \mathbb{I} \otimes \mathbb{I} \otimes \mathbb{I} + \gamma^0 \otimes \gamma^0 \otimes \mathbb{I} + \gamma^0 \otimes \mathbb{I} \otimes \gamma^0 + \mathbb{I} \otimes \gamma^0 \otimes \gamma^0 \right]$$

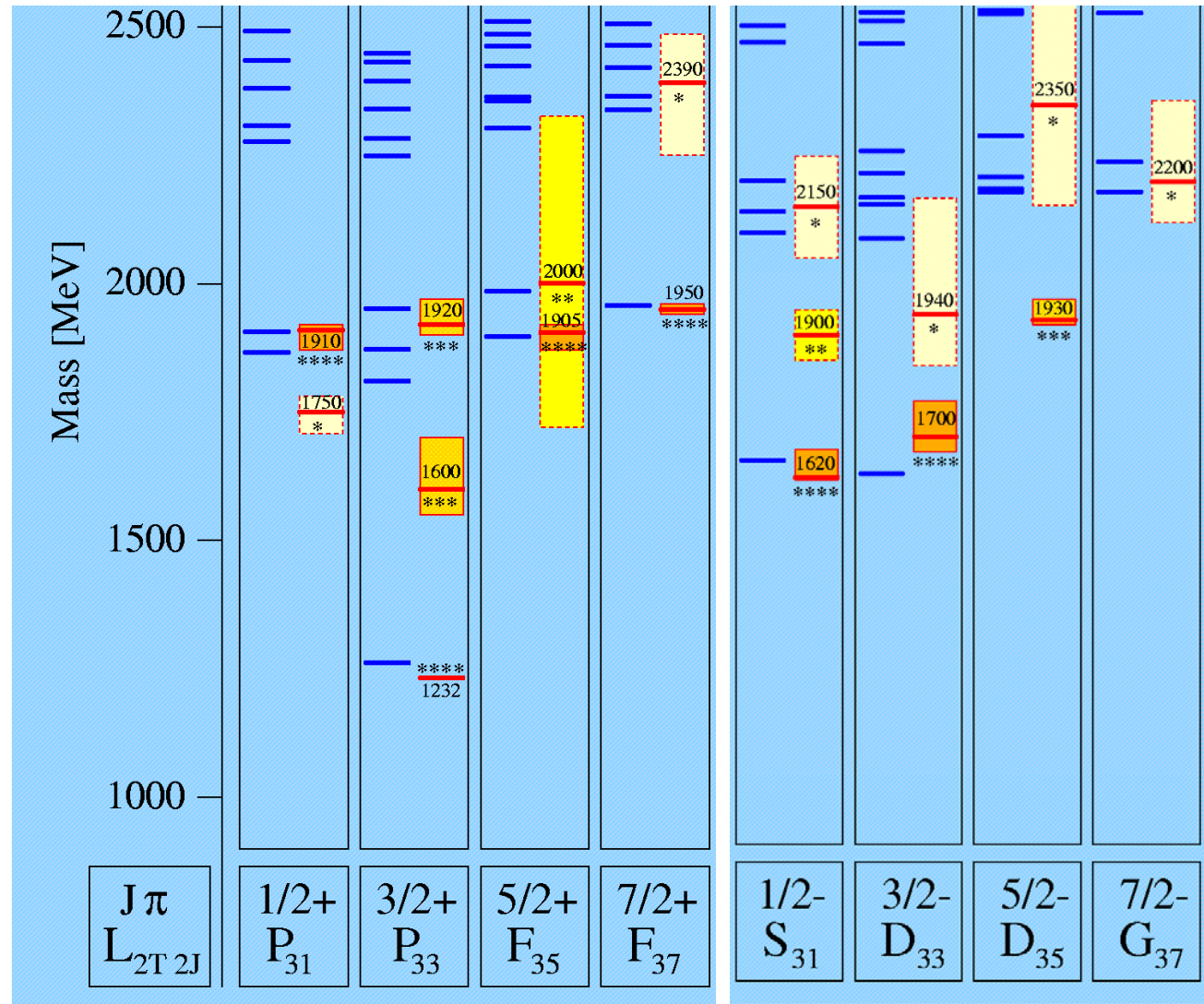
$$B_3 = b \frac{1}{2} \left[ -\mathbb{I} \otimes \mathbb{I} \otimes \mathbb{I} + \gamma^0 \otimes \gamma^0 \otimes \mathbb{I} + \gamma^0 \otimes \mathbb{I} \otimes \gamma^0 + \mathbb{I} \otimes \gamma^0 \otimes \gamma^0 \right]$$

- Chosen to reduce spin-orbit effects
- Produces correct Regge trajectories
- For  $N^*, D^*$  only five parameters:  
 $m_n, g_{nn}, a, b, \lambda$



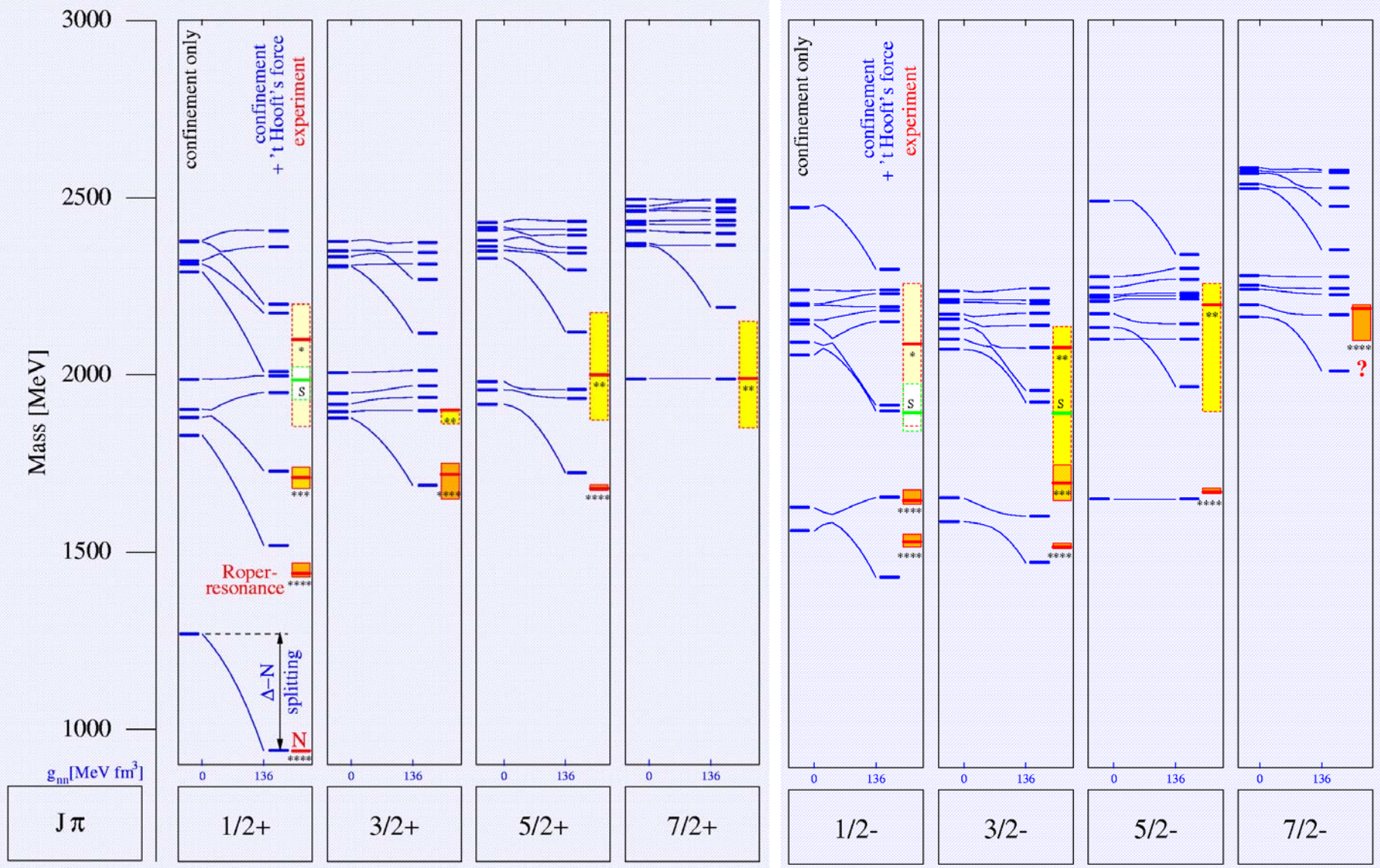
# Instanton-induced interactions...

- spectrum of  $\Delta^*$  only from confining potential
- Blask, Bohn, Huber, Metsch & Petry





# $N^*$ spectrum from 't Hooft's force



# Nature of short-range qq interactions

- Spectrum alone cannot distinguish models
  - Wave functions much more sensitive
- Baryons have near degenerate excited states
  - two orbital spaces in 3q system, degenerate  $l=1$  orbital excitations
  - two possible total quark spins
  - Degeneracy broken by spin-spin interactions ( $L=0 \times S=0$  operator)
  - Large mixing caused by tensor, spin-orbit,... interactions which must be present (relativity)





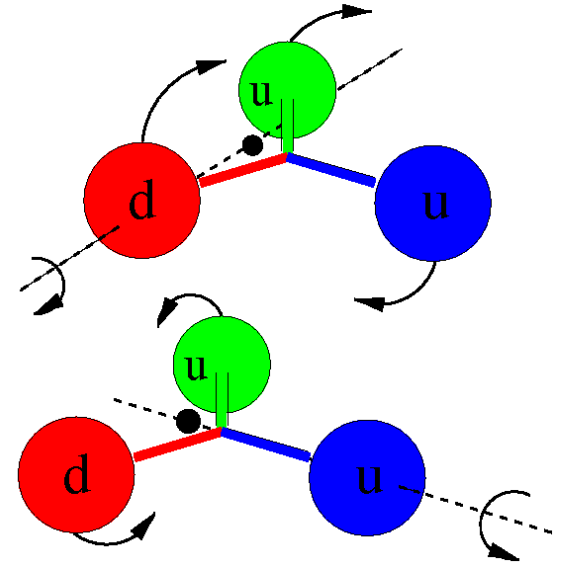
# Why are mixing angles important?...

- Chizma and Karl, hep-ph/0210126
  - Different effective interactions, adjusted to contact splittings, have very different tensor (and spin-orbit) interactions
    - Cause different mixing angles in low-lying  $S_{11}$  &  $D_{13}$  states
- ⇒ A determination of mixing angles can establish nature of effective interactions between quarks



# Low lying excited baryons

- Negative-parity excited states
$$L^P=1^- + S_{qqq}(1/2) = 1/2^-, 3/2^-$$
$$L^P=1^- + S_{qqq}(3/2) = 1/2^-, 3/2^-, 5/2^-$$
- Orbital excitations in  $\rho$  or  $\lambda$
- $2(N1/2^-, N3/2^-), N5/2^-, D1/2^-, D3/2^-$ 
  - Two  $N1/2^-$  model states,  
 $S=1/2$  or  $S=3/2$ ,
    - are  $S_{11}$  in  $\pi N \Rightarrow \pi N$
    - resonances  $N(1535)$  and  $N(1650)$  in analyses
  - Two  $N3/2^-$  states are  $D_{13} \Rightarrow N(1520), N(1700)$



# Low lying excited baryons...

- Without spin-dependent interactions between the quarks
  - Two degenerate  $S_{11}$  states ( $S=1/2$  and  $S=3/2$ ) and two degenerate  $D_{13}$  states
- Contact interactions breaks degeneracy
  - Proportional to  $\sum_{i<j} \mathbf{S}_i \cdot \mathbf{S}_j f(r_{ij})$
  - Tensor and spin-orbit interactions mix two  $S=1/2$  and  $S=3/2$  states with same  $J^P$



# Mixing angles

- Physical states are admixtures of two possible L,S combinations

$$N(1535)1/2^- = \cos(\theta_S) N^2P1/2^- - \sin(\theta_S) N^4P1/2^-$$

$$N(1650)1/2^- = \sin(\theta_S) N^2P1/2^- + \cos(\theta_S) N^4P1/2^-$$

$$N(1520)3/2^- = \cos(\theta_D) N^2P3/2^- - \sin(\theta_D) N^4P3/2^-$$

$$N(1700)3/2^- = \sin(\theta_D) N^2P3/2^- + \cos(\theta_D) N^4P3/2^-$$

- How can  $\theta_S$  and  $\theta_D$  be determined?
- Lattice QCD, with enough time (CPU and elapsed!) and clever choice of correlators... or model of baryon structure and strong decays



# Unmixed model states

- Work done with **W. Roberts (Jlab/ODU)**
- Variational calculation in large HO basis
  - Unmixed states are orthogonal eigenstates of  $H$ , with definite  $L(=1)$  and  $S_{qq}(1/2 \text{ or } 3/2)$
- OGE contact and color-Coulomb interactions
  - Turn off any interactions which are not simultaneous spin and spatial scalars
  - No tensor, spin-orbit (from confinement, OGE)

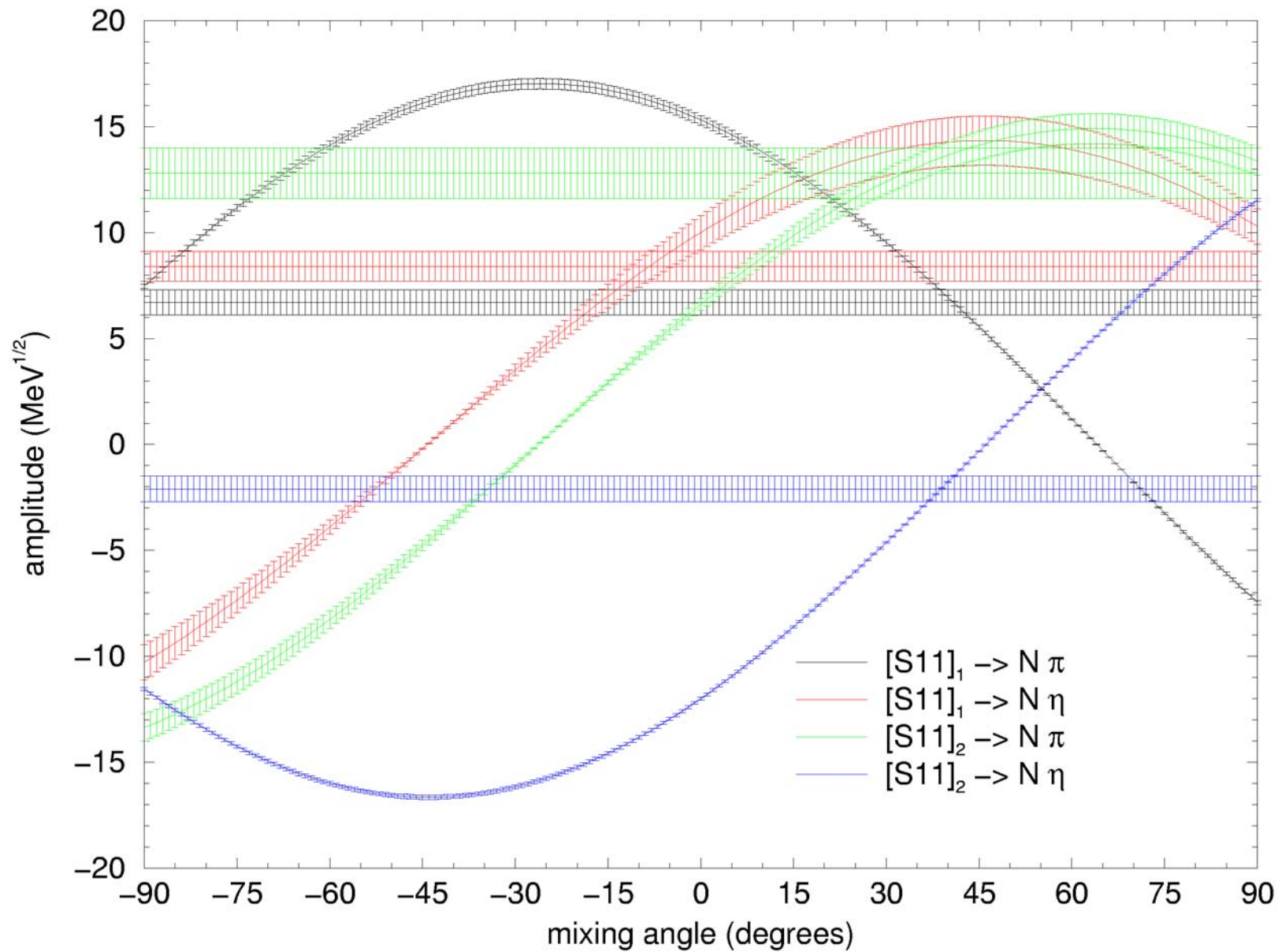


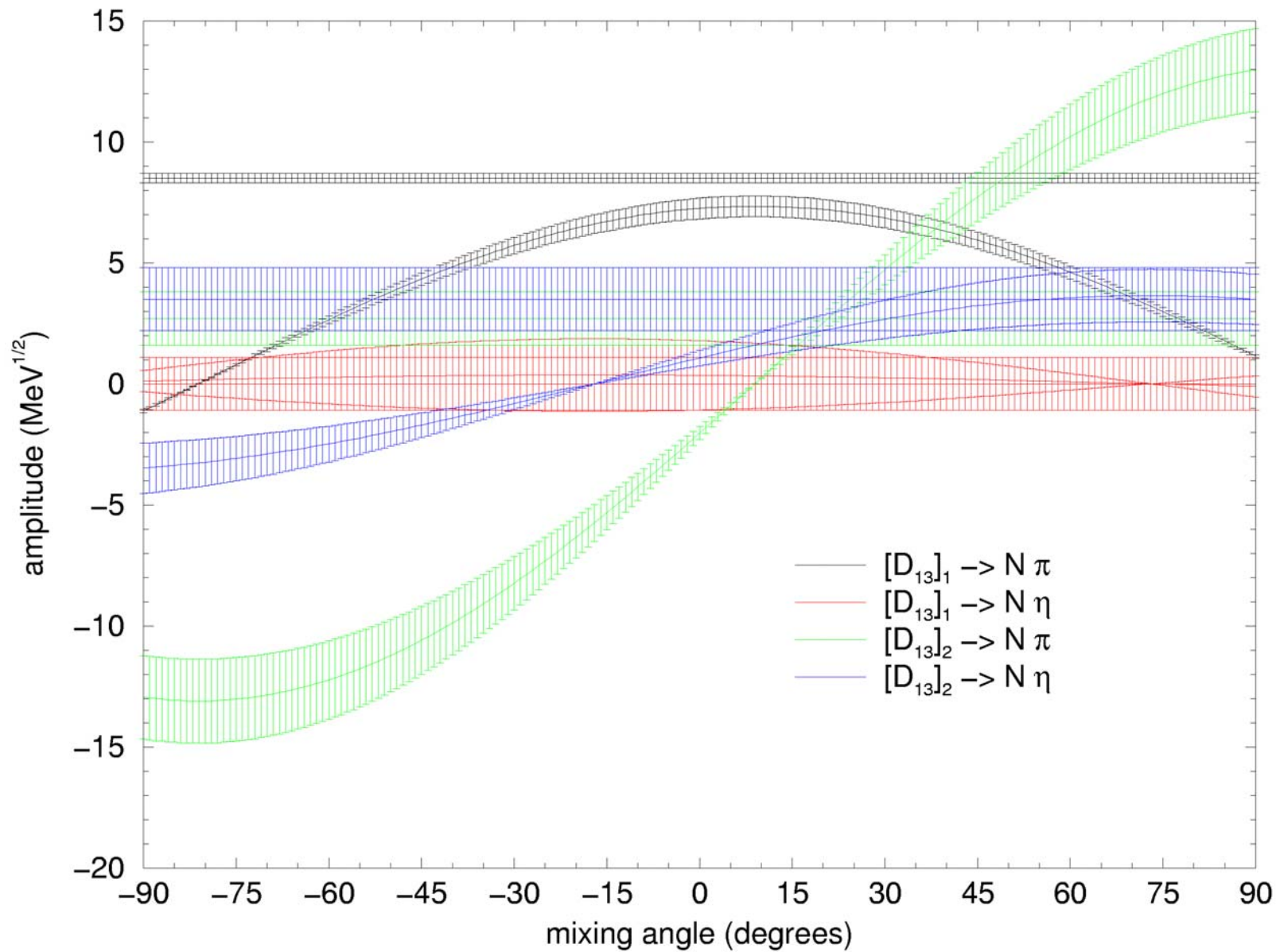
# Decay model

- Use  $^3P_0$ , phenomenological decay model
  - Calculate decay amplitudes of unmixed states to  $N\pi$ ,  $N\eta$ ,  $\Lambda K$
- Find decay amplitudes as a function of  $\theta_S$  and  $\theta_D$
- Compare to Pitt-ANL analysis results









# Extracted mixing angles

	$[ ]_1 \Rightarrow \pi N$	$[ ]_2 \Rightarrow \pi N$	$[ ]_1 \Rightarrow \eta N$	$[ ]_2 \Rightarrow \eta N$	$[ ]_2 \Rightarrow \Lambda K$
$\theta_S$	$41^\circ$	$33^\circ$	$-8^\circ$	$39^\circ$	$\sim 20^\circ$
$\theta_D$	$9^\circ$	$21^\circ$		$> 11^\circ$	

- Model does not fully describe  $S_{11}(1535) \Rightarrow \eta N$
- Mixing angles  $\theta_S \sim 35-40^\circ$ ,  $\theta_D \sim 10-20^\circ$
- Other determinations:  $\theta_S \sim 32^\circ$ ,  $\theta_D \sim -10^\circ$ 
  - B. Saghai and Z. Li
  - Hey, Litchfield and Cashmore ('75)



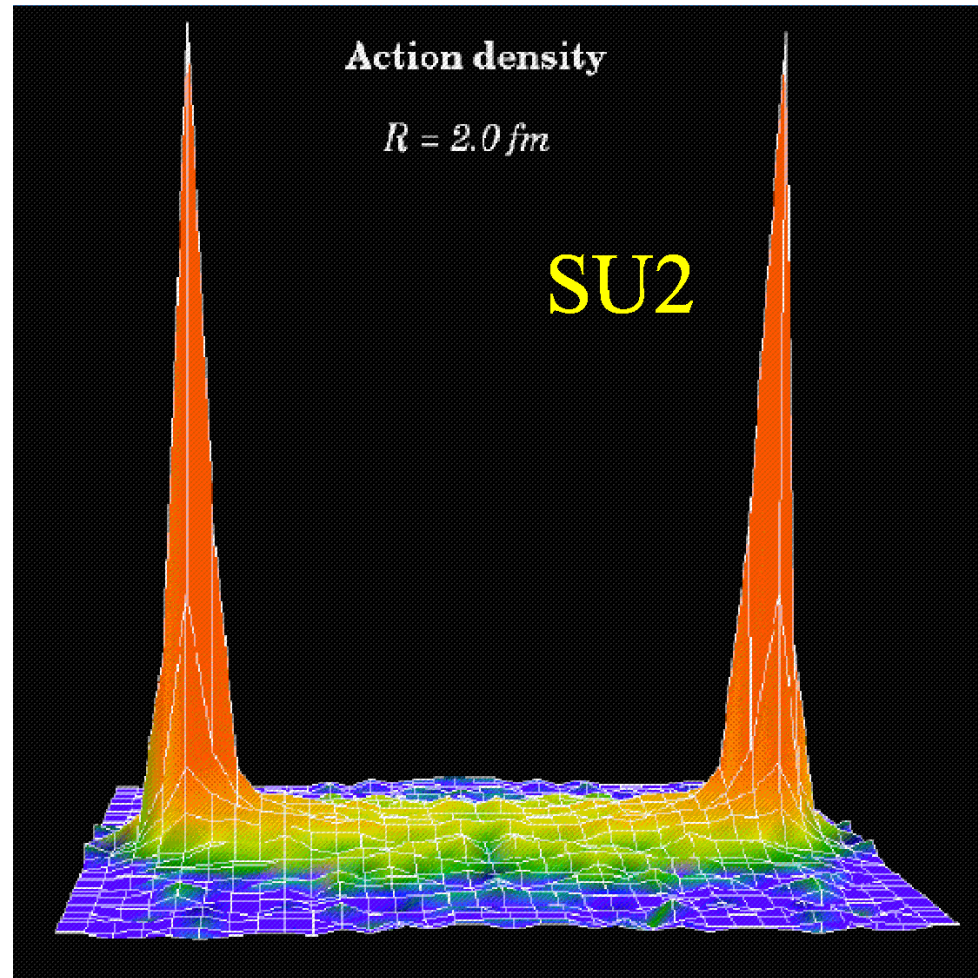
# New insights from lattice QCD

- Is the potential model picture reasonable, at least for heavy-quark systems?
- Can lattice calculations explain the position of the light positive-parity states (like the Roper)?



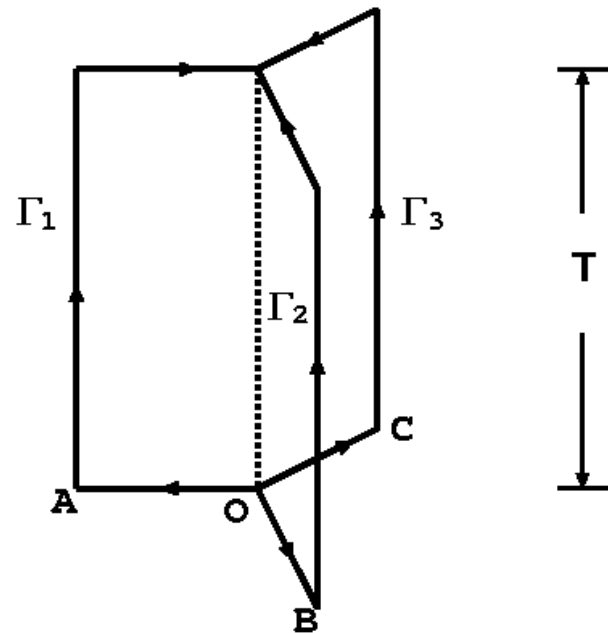
# Confinement in heavy quark mesons

- Bali et al.
- Static quark and antiquark separated by 2 fm
- Gluonic flux shows tube structure between quarks



# How should we treat confinement?

- Quenched lattice measurement of QQQ potential
- Takahashi, Matsufuru, Nemoto and Suganuma, PRL86 (2001) 18.
- Measure potential with 3Q-Wilson loop (static quarks) for  $0 < t < T$
- Also fit  $Q\bar{Q}$  potential to compare  $\sigma$  and Coulomb terms





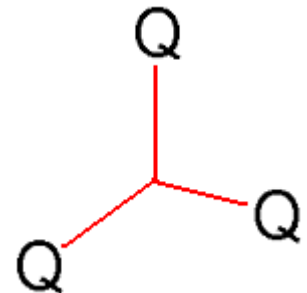
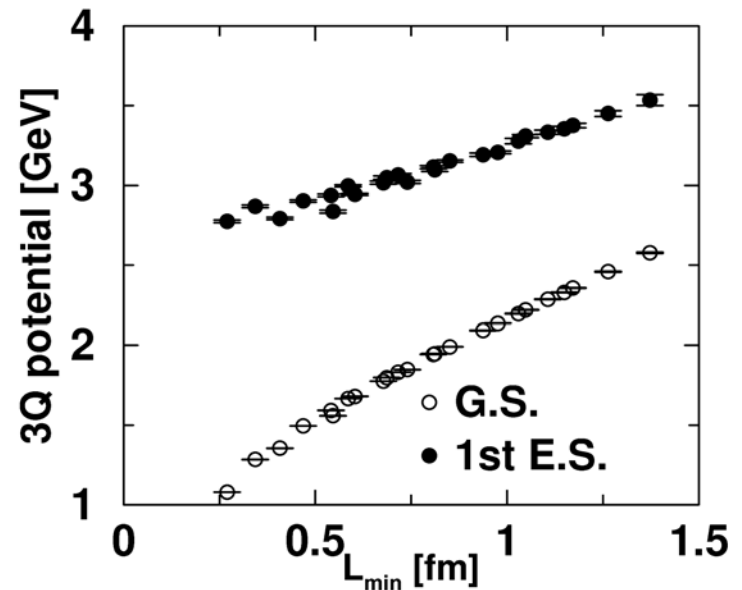
# Confinement

- Quenched lattice calculation of QQQ static potential

- Takahashi, Matsufuru, Nemoto and Suganuma, PRL **86** (2001) 18, **hep-lat/0210024**
- Plot vs.  $L_{\min}$  = min. length Y-shaped string for 24 configurations

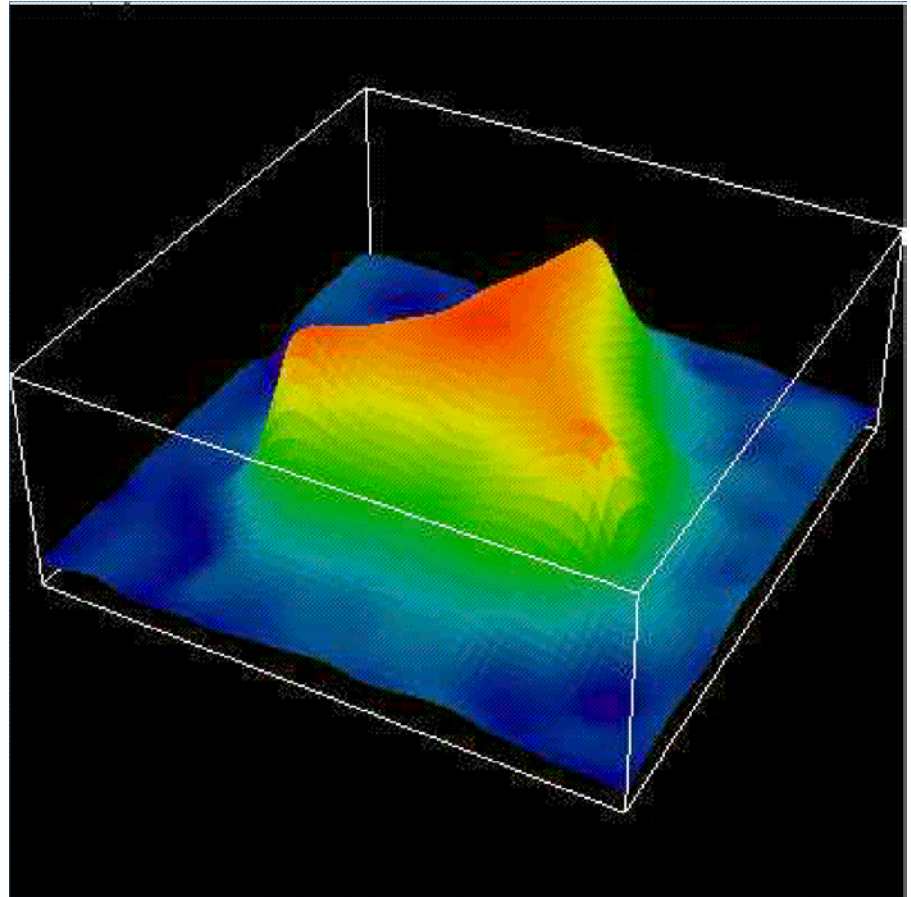
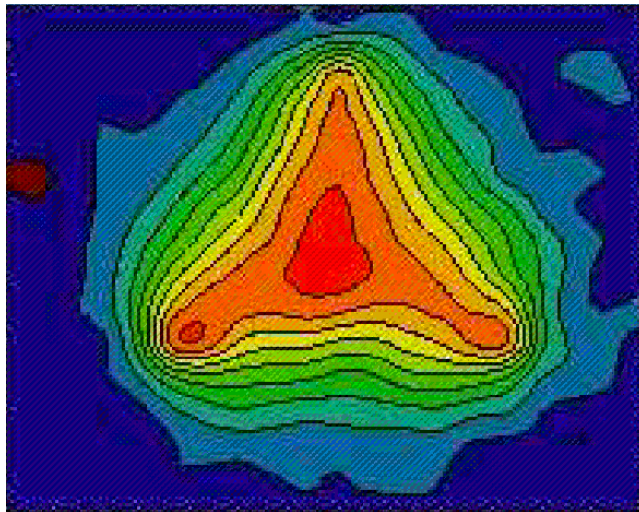
⇒ String-like flux-tube potential in QQQ baryons

$$V_B(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \sigma(l_1 + l_2 + l_3) = \sigma L_{\min}$$



# Flux tube in QQQ baryons

- Abelian action distribution of gluons and light quarks nr. QQQ
- Ichie, Bornyakov, Struer & Schierholz, hep-lat/0212024



# Leading Born-Oppenheimer approximation

- Juge, Kuti, Morningstar PRL82, 4400 ('99)
  - Born Oppenheimer approximation for heavy-quark mesons
  - Use heavy-quark adiabatic potentials from lattice, including excited glue states
  - No light-quarks, no quark spin or retardation effects
- Compared to NRQCD results (Morningstar)
  - Retardation effects and mixing between states in different adiabatic potentials allowed



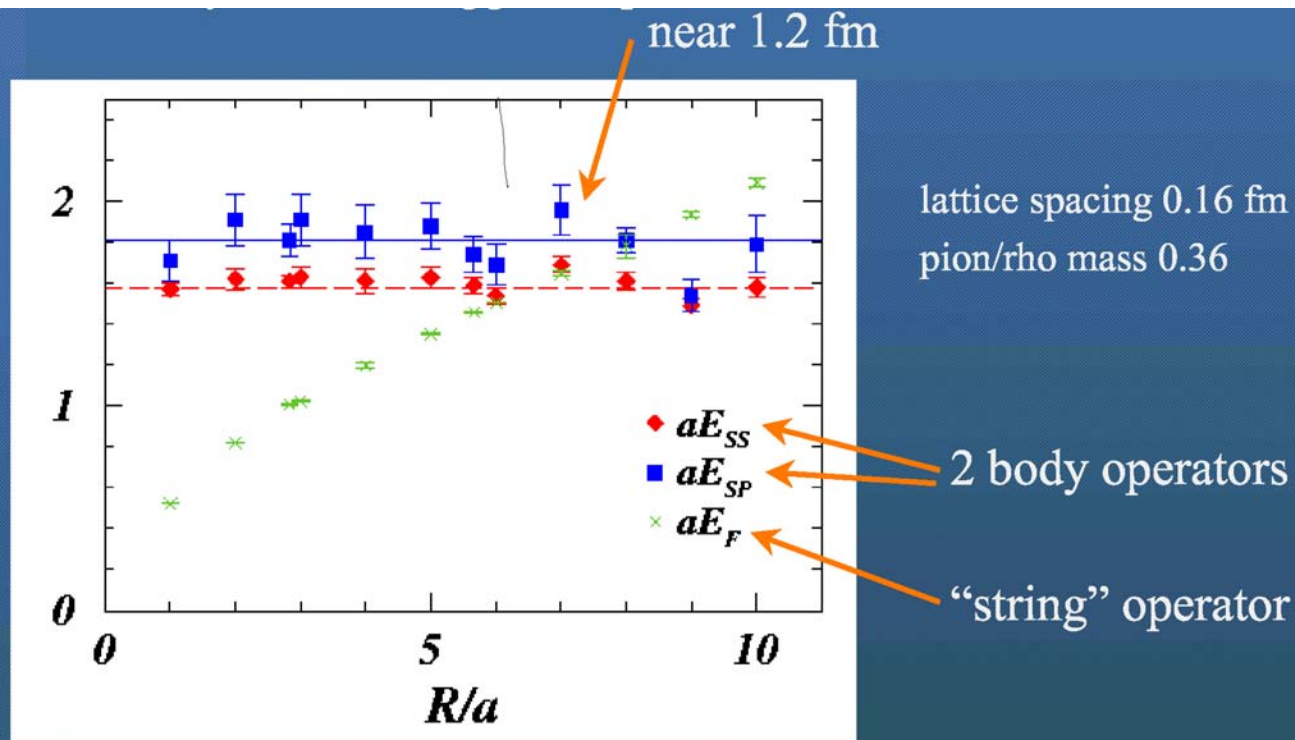
# LBO approximation...

- Good agreement (10%) for level splittings of 4 conventional and 2 hybrid heavy-quark mesons
  - Partially explains success of constituent quark model
  - Allows inclusion of gluon dynamics
- Current LQCD studies:
  - Does LBO survive inclusion of quark spin?
  - OK, but worse in charmonium than upilon
    - Drummond et al PLB478, 151 ('00)
    - Burch and Toussaint, hep-lat/0305008
  - Do light dynamical quarks spoil this picture?



# Light quark effects on potential?

- Explore string breaking using two-body operators
- Allow two flavors of dynamical staggered quarks



Bernard et al., PRD64, 074509 (2001)



# Baryon spectroscopy in lattice QCD

- Recent progress in understanding first excited states ( $J^\pi=1/2^+, 1/2^-$ )
  - If pion masses  $\sim 500$  MeV, calculated masses  $N_{1/2^+} < N_{1/2^-} < N^*_{1/2^+}$  with roughly equal spacing
    - Physical state at (940/1535/1440)
  - Problem with quenched approximation?
  - Problem with  $m_\pi$  (light quark masses) being too large?



# Baryon spectroscopy in LQCD...

- C. Maynard, D.G. Richards (LHPC/UKQCD)
  - Show that this remains true in full (unquenched) QCD at  $m_\pi \sim 500$  MeV
- Light quarks ( $m_\pi$ ):
  - Wilson fermions:  $m_\pi \sim 400$  MeV
  - Smoothed actions (FLIC)  $m_\pi \sim 300$  MeV
    - Have chiral symmetry only in  $a \rightarrow 0$  limit
- Overlap/domain wall fermions
  - Have exact analog of  $\chi$ -symm. at finite  $a$
  - Other physics and calculational advantages
  - But cost a factor of 30 in CPU time!

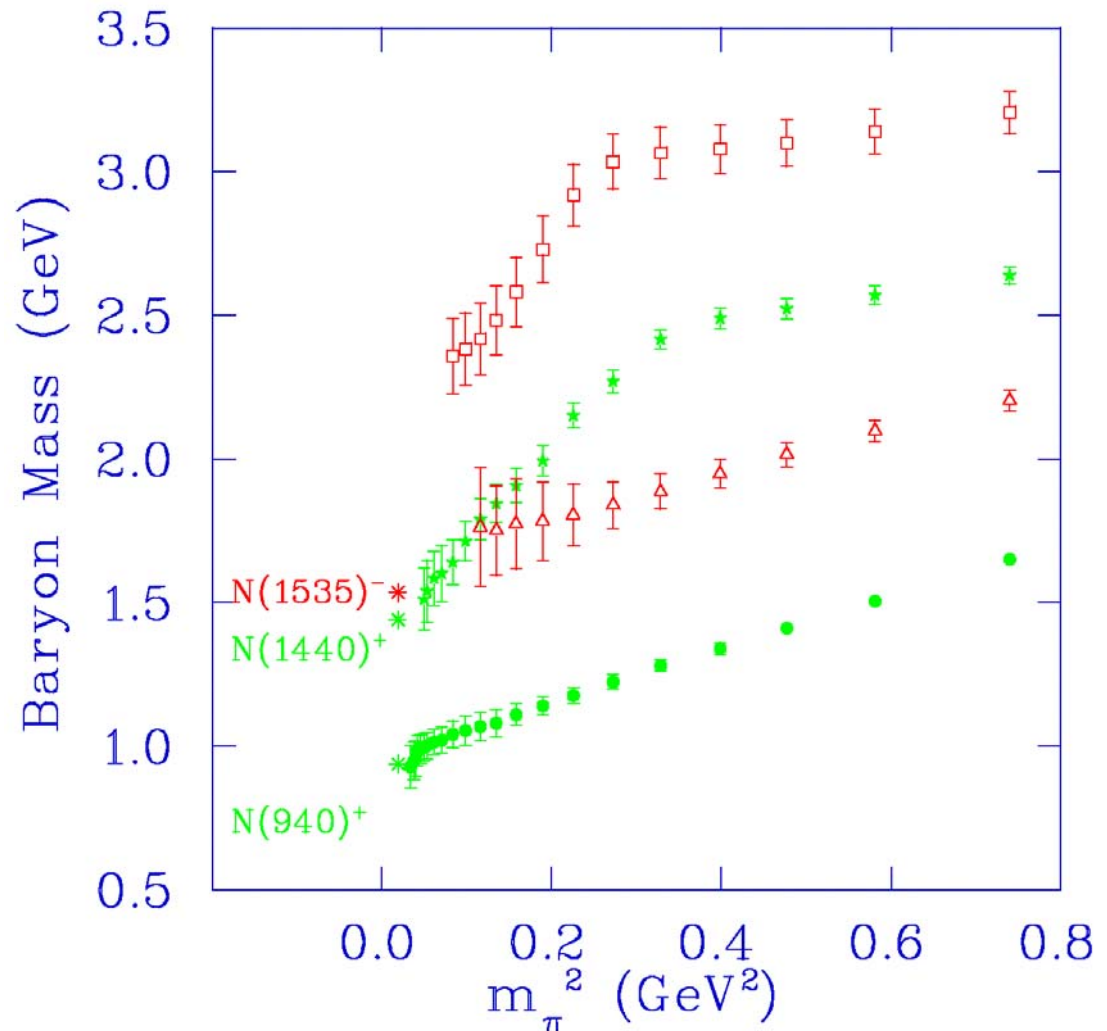




# Light $m_\pi$ baryon masses in LQCD

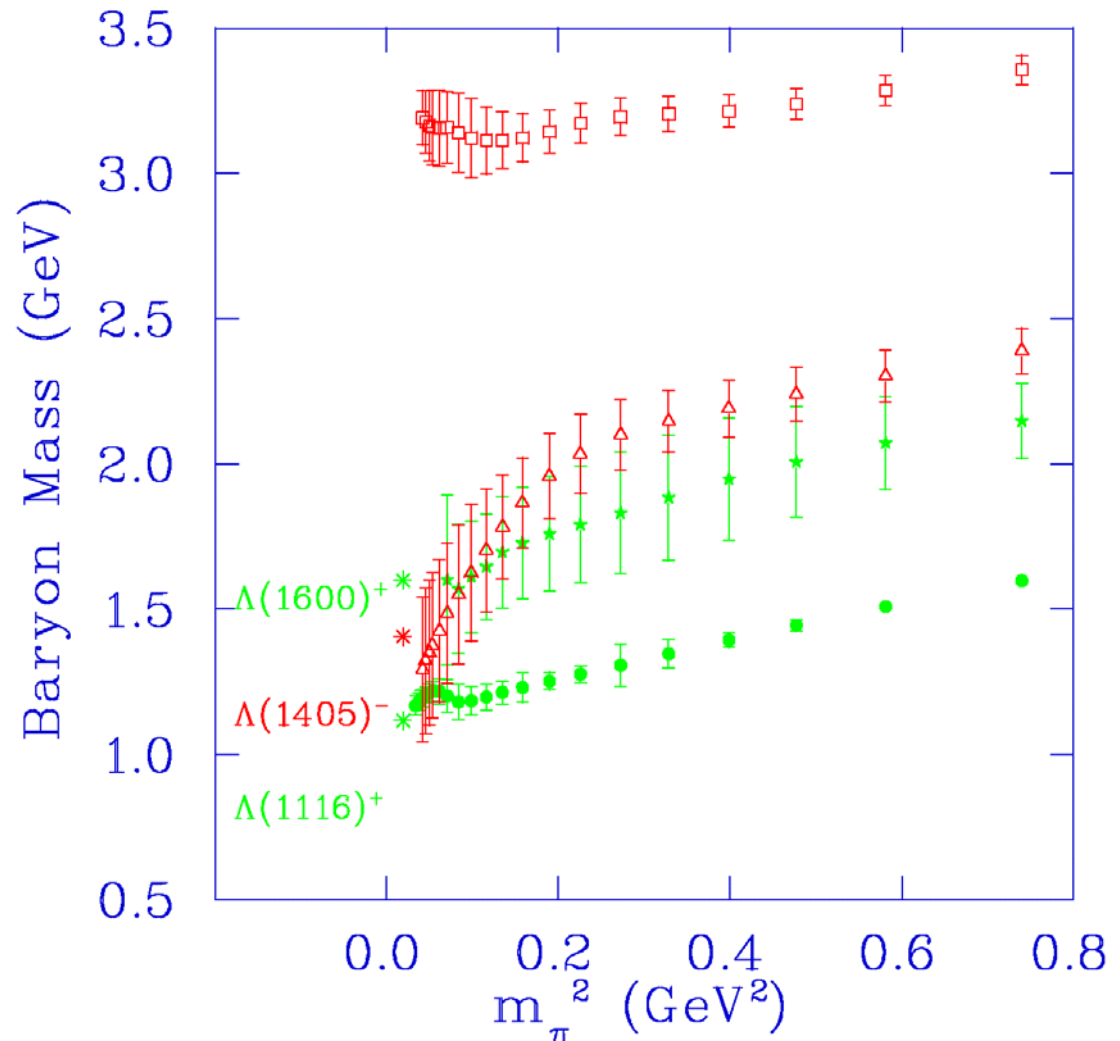
F. Lee et al.  
(Kentucky  
collaboration)  
hep-lat/0208070

- Use overlap fermions and simple  $qqq$  operators
- Levels cross at low  $m_\pi$



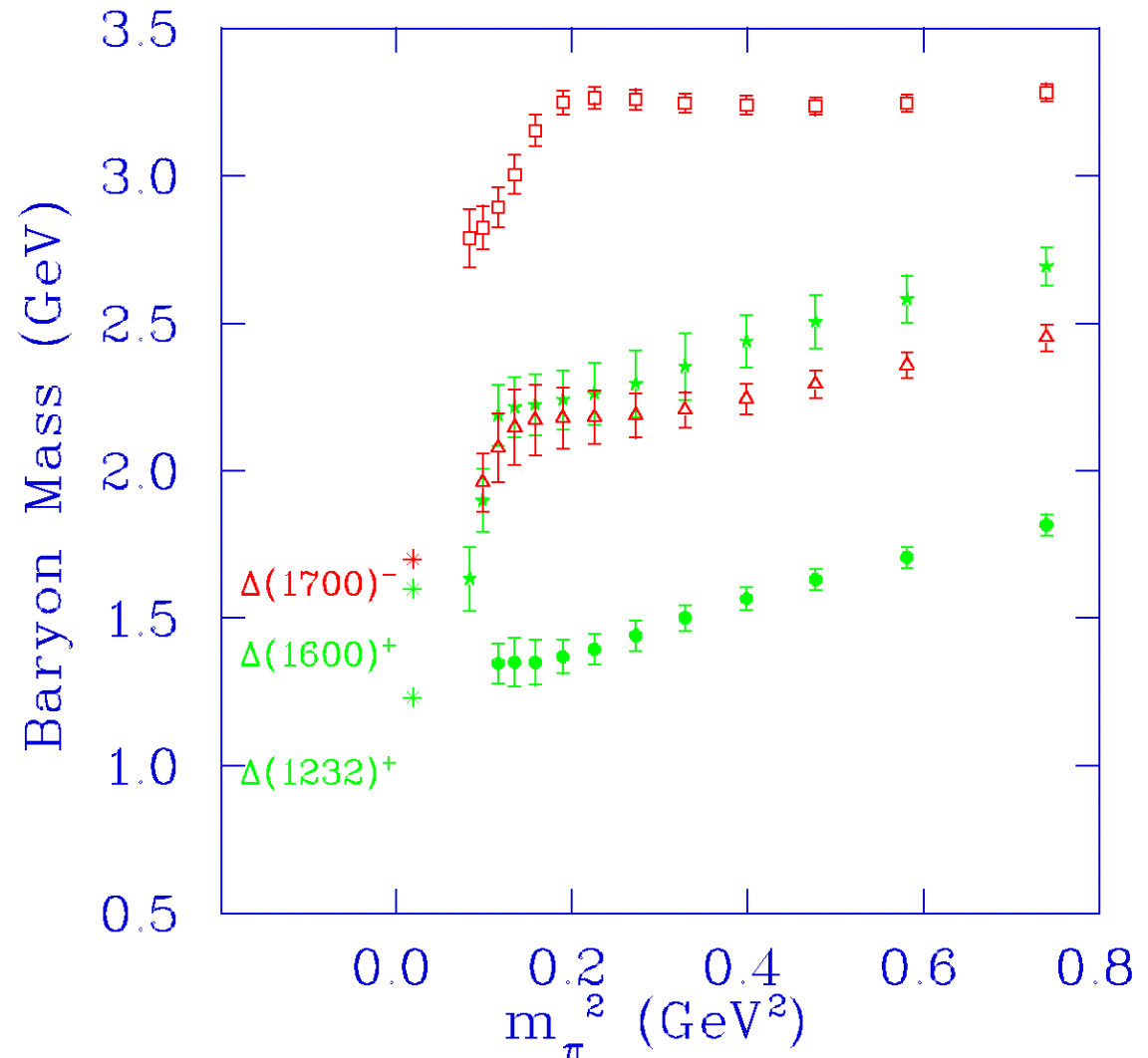
# Light $m_\pi$ baryon masses in LQCD...

- Same level crossing observed in light  $\Lambda 1/2^+$  and  $\Lambda 1/2^-$  states



# Light $m_\pi$ baryon masses in LQCD...

- Also shows approximate degeneracy of  $\Delta^{*3/2+}$  and  $\Delta^{3/2-}$



# Baryon spectroscopy in LQCD...

- Chiral behavior is important!
  - Non-analytic terms present in low  $m_\pi$  limit give rapidly varying behavior
  - Need to make (quenched)  $\chi$ -PT extrapolations of lattice data to low  $m_\pi$
- Roper likely a  $qqq$  state
  - Bare quark model  $qqq$  state should be too massive!
  - Couples strongly to  $N\pi$  decay channel
  - Model with baryon-meson loops?
- Constituent quark model without loop effects will miss this important physics



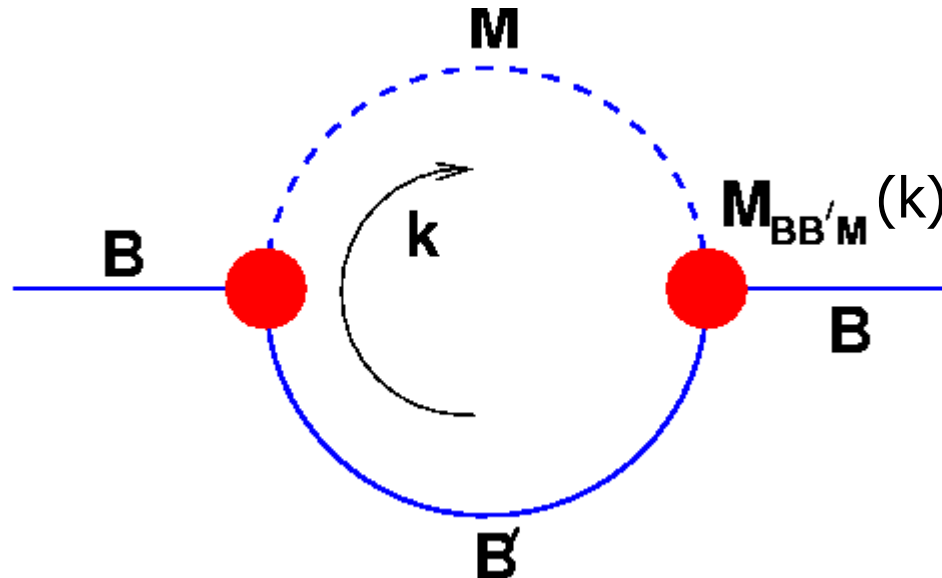
# Loop effects in baryon spectrum

- Two approaches:
  - Include elementary (bare)  $qqq$  excitations in careful calculation of reaction observables
    - Include rescattering into open  $B'M$  channels
    - S. Krewald, et al. (Juelich), T. Sato & H. Lee, C. Bennhold & GWU group
  - Explicitly evaluate hadron loop corrections to masses and decays
    - Mesons ( $\rho$ - $\omega$ ): P. Geiger and N. Isgur



# Unquenching the quark model

- In QCD  $qqq(q\bar{q})$  configurations possible in baryons: effect on CQM?
  - Model with baryon-meson intermediate states, loops  $\Rightarrow$  self energies: **Danielle Morel & SC**
  - High-momentum part of loops contains OBE



# Unquenching the quark model...

- Hecht, Roberts, Tandy, Thomas,...
  - Schwinger-Dyson Bethe-Salpeter study
  - examine  $m_\pi$  dependence of  $N\pi$  loop contribution to nucleon mass
- D. Morel and A. Thomas
  - Studying  $m_\pi$  dependence of contributions to resonance masses from  $N\pi$  &  $\Delta\pi$ ,... loops
  - Non-analytic behavior in extrapolation of lattice data for baryon masses to light quark (pion) masses





# Summary

- New data from JLab and elsewhere
  - Contains evidence for new (missing) baryon states
  - Refines our knowledge of existing states
- Unitary multi-channel analysis of this data
  - Will establish new states
  - Can distinguish between QCD-inspired models
- Dramatic new lattice calculations
  - Support flux-tube potential model picture, at least for heavy-quark hadrons
  - Point to importance of decay-channel couplings with light quarks



# Summary...

- Challenges to theory:
  - Analysis of data in many channels:
    - Reduce ambiguities in analysis by incorporating constraints from:
      - Multi-channel unitarity, analyticity, gauge invariance,...
    - Efficient treatment of 'background' consistent with these constraints
  - Include effects of open channels in:
    - Lattice QCD calculations
    - Reaction models
    - QCD-inspired models





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Jlab Users Group Mtg. 2003-55